# A $C^{0}$ finite element method for the biharmonic problem with Navier boundary conditions in a polygonal domain 

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In this paper we study the biharmonic equation with Navier boundary conditions in a polygonal domain. In particular, we propose a method that effectively decouples the fourth-order problem as a system of Poisson equations. Our method differs from the naive mixed method that leads to two Poisson problems but only applies to convex domains; our decomposition involves a third Poisson equation to confine the solution in the correct function space, and therefore can be used in both convex and nonconvex domains. A $C^{0}$ finite element algorithm is in turn proposed to solve the resulting system. In addition, we derive optimal error estimates for the numerical solution on both quasi-uniform meshes and graded meshes. Numerical test results are presented to justify the theoretical findings.

Keywords: biharmonic equation; reentrant corner; mixed formulation; $C^{0}$ finite element method; optimal error estimates.

## 1. Introduction

Let $\Omega \subset \mathbb{R}^{2}$ be a polygonal domain. Consider the biharmonic problem

$$
\begin{equation*}
\Delta^{2} u=f \quad \text { in } \Omega, \quad u=0 \quad \text { and } \quad \Delta u=0 \quad \text { on } \partial \Omega \tag{1.1}
\end{equation*}
$$

[^0]The boundary conditions in (1.1) are referred to as the homogeneous Navier boundary conditions (Destuynder \& Salaun, 1996; Rafetseder \& Zulehner, 2018) that occur for example in the model for the static loading of a pure hinged thin plate. Equation (1.1) is a fourth-order elliptic equation for which a direct finite element approximation usually involves delicate construction of the finite element space (Argyris et al., 1968; Brenner \& Scott, 2002; Ciarlet Jr \& He, 2003). An alternative approach is to use a mixed formulation to decompose the high-order problem into a system of equations that may be easier to solve. This approach is particularly appealing for the biharmonic problem (1.1) because the Navier boundary condition allows one to obtain two Poisson equations that are completely decoupled, which implies that a reasonable numerical solution should be achieved by merely applying a finite element Poisson solver to the mixed formulation. However, it has been observed (Nazarov \& Sweers, 2007a; Zhang \& Zhang, 2008; Gerasimov et al., 2012; De Coster et al., 2019) that the performance of this standard mixed method depends on the domain geometry. In a convex domain, the corresponding numerical approximations converge to the solution of equation (1.1), although the convergence rate may not be optimal. When the domain possesses reentrant corners, however, the result can be misleading: this mixed finite element formulation produces numerical solutions that may be converging to a wrong solution.

In this paper we propose a $C^{0}$ finite element algorithm for solving the biharmonic problem (1.1) and analyze its efficiency. In particular, we shall devise a modified mixed formulation to transform equation (1.1) into a system of three Poisson equations. This work is based on the observation that the aforementioned standard mixed formulation (decomposition into two Poisson equations) in fact defines a weak solution in a space larger than that for equation (1.1). This mismatch in function spaces does not affect the solution in a convex domain, while in a nonconvex domain it allows additional singular functions and therefore results in a solution different from that of equation (1.1) (known as the Sapongyan paradox; Nazarov \& Plamenevsky, 1994; Zhang \& Zhang, 2008). Our modified mixed formulation ensures that the associated solution is identical to the solution of (1.1) in both convex and nonconvex domains. This is accomplished by introducing an additional intermediate Poisson problem that confines the solution in the correct space.

To solve the proposed mixed formulation we present a numerical algorithm based on the piecewise linear $C^{0}$ finite element. Moreover, we derive the error analysis on the finite element approximations for both the auxiliary function $w$ (see (2.11)) and the solution $u$. For the auxiliary function $w$, the error in the $H^{1}$ norm is standard and has a convergence rate $h^{\frac{\pi}{\omega}}$ on a quasi-uniform mesh, where $\omega$ is the interior angle of the reentrant corner; its $L^{2}$ error estimate can be obtained using a duality argument. For the solution $u$, the error in the $H^{1}$ norm is bounded by (i) the interpolation error of the solution $u$ in $H^{1}$, (ii) the $L^{2}$ error for the auxiliary function $w$ and (iii) the $L^{2}$ error for the solution $\xi$ of the additional intermediate Poisson problem. We shall show that the proposed algorithm has the optimal $H^{1}$ convergence rate for the solution $u$ on quasi-uniform meshes.

In addition, we present regularity estimates for the proposed system in a class of Kondratievtype weighted spaces. Based on these regularity results, we in turn propose graded mesh refinement algorithms, such that the associated finite element methods recover the optimal convergence rate in the energy norm for the auxiliary function $w$ even when $w$ is singular. For clarity in the exposition, we adopt the linear $C^{0}$ finite element method in this paper with the assumption that the domain $\Omega$ has at most one reentrant corner. The cases involving high-order finite elements and multiple reentrant corners will be discussed in a forthcoming paper.

There is a rich literature on mixed finite element methods for the biharmonic problem, and most existing works are focused on the Dirichlet (clamped) boundary condition ( $u=\partial_{n} u=0$ on $\partial \Omega$ ). For example, Davini \& Pitacco (2000) studied a mixed method, and in Zulehner (2015)
preconditioning techniques were investigated for the discrete system, both of which are based on the Ciarlet-Raviart formulation (Ciarlet \& Raviart, 1974) for solving biharmonic problems with a clamped boundary condition. In addition, some other types of mixed variational formulations can be found in Gallistl (2017), Sweers (2009) for the clamped biharmonic problem. Notice that it is possible to extend the algorithm in Gallistl (2017) to solve equation (1.1). We also mention the related works Blum et al. (1980), De Coster et al. (2019) and the references therein for general regularity results for biharmonic problems, Bernardi et al. (1992) for a mixed spectral element method for the Navier-Stokes equations and Nazarov \& Sweers (2007b) for Kirchhoff plate bending problems. In particular, De Coster et al. (2019) studied the Saponzhyan-Babushka paradox of the biharmonic problem with piecewise $C^{2,1}$ boundaries connecting at corners, and an augmented solution space is needed to correct the solution of the second-order system in the appropriate Sobolev-type space.

The method proposed in this paper effectively decouples the biharmonic problem (1.1) into a system of Poisson equations, whose solution is equivalent to that of the original problem in both convex and nonconvex domains. The $C^{0}$ finite element algorithm is simple, robust and effective in practical computations for problem (1.1). Our method can be further applied to cases involving multiple reentrant corners and high-order finite element approximations. Moreover, with suitable modifications, we expect to extend this algorithm to some other high-order problems, such as a class of fourth-order problems with low-order terms and sixth-order problems with similar boundary conditions. It is also reasonable to explore the extension of this work to fourth-order problems with nonlinear operators or other boundary conditions.

The rest of the paper is organized as follows. In Section 2, based on the general regularity theory for second-order elliptic equations (Kondrat'ev, 1967; Moussaoui, 1985; Grisvard, 1992; Kozlov et al., 1997, 2001), especially Nazarov \& Svirs (2007a), we review the weak solutions of the biharmonic problem (1.1) and the naive mixed formulation. In addition, we discuss the orthogonal space of the image of the operator $-\Delta$, which is one-dimensional, and identify a basis function in this space. Then we propose a modified mixed formulation and show the equivalence of its solution to the original biharmonic problem. In Section 3 we propose the finite element algorithm and obtain error estimates on quasi-uniform meshes for both the solution $u$ and the auxiliary function $w$. In Section 4 we introduce a weighted Sobolev space and derive regularity estimates for the solution near the reentrant corner. Then we present the graded mesh algorithm and provide optimal error estimates on graded meshes. We report numerical test results in Section 5 to validate the theory.

Throughout the paper, the generic constant $C>0$ in our estimates may be different at different occurrences. It will depend on the computational domain, but not on the functions involved nor on the mesh level in the finite element algorithms.

## 2. The biharmonic problem

### 2.1 Well-posedness of the solution

Denote by $H^{m}(\Omega), m \geq 0$ the Sobolev space that consists of functions whose $i$ th derivatives are square integrable for $0 \leq i \leq m$. Let $L^{2}(\Omega):=H^{0}(\Omega)$. Recall that $H_{0}^{1}(\Omega) \subset H^{1}(\Omega)$ is the subspace consisting of functions with zero trace on the boundary $\partial \Omega$. We shall denote the norm $\|\cdot\|_{L^{2}(\Omega)}$ by $\|\cdot\|$ when there is no ambiguity about the underlying domain. For $s>0$, let $s=m+t$, where $m \in \mathbb{Z}_{\geq 0}$ and $0<t<1$. Recall that for $D \subseteq \mathbb{R}^{d}$, the fractional-order Sobolev space $H^{s}(D)$ consists of distributions $v$
in $D$ satisfying

$$
\|v\|_{H^{s}(D)}^{2}:=\|v\|_{H^{m}(D)}^{2}+\sum_{|\alpha|=m} \int_{D} \int_{D} \frac{\left|\partial^{\alpha} v(x)-\partial^{\alpha} v(y)\right|^{2}}{|x-y|^{d+2 t}} \mathrm{~d} x \mathrm{~d} y<\infty,
$$

where $\alpha=\left(\alpha_{1}, \ldots, \alpha_{d}\right) \in \mathbb{Z}_{\geq 0}^{d}$ is a multi-index such that $\partial^{\alpha}=\partial_{x_{1}}^{\alpha_{1}} \cdots \partial_{x_{d}}^{\alpha_{d}}$ and $|\alpha|=\sum_{i=1}^{d} \alpha_{i}$. Let $\widetilde{H}^{s}(D)$ be the space of all $v$ defined in $D$ such that $\tilde{v} \in H^{s}\left(\mathbb{R}^{d}\right)$, where $\tilde{v}$ is the extension of $v$ by zero outside $D$.

The following variational formulation for equation (1.1) can be obtained using integration by parts:

$$
\begin{equation*}
a(u, v):=\int_{\Omega} \Delta u \Delta v \mathrm{~d} x=\int_{\Omega} f v \mathrm{~d} x=(f, v) \quad \forall v \in H^{2}(\Omega) \cap H_{0}^{1}(\Omega) . \tag{2.1}
\end{equation*}
$$

For a function $v \in H^{2}(\Omega) \cap H_{0}^{1}(\Omega)$, applying the Poincaré-type inequality (Grisvard, 1992, Theorem 2.2.3) gives $\|\Delta v\|_{L^{2}(\Omega)} \geq C\|v\|_{H^{2}(\Omega)}$. Then by the Lax-Milgram theorem, equation (2.1) defines a unique weak solution $u \in H^{2}(\Omega) \cap H_{0}^{1}(\Omega)$ for any $f$ in the dual space of $H^{2}(\Omega) \cap H_{0}^{1}(\Omega)$ (namely, $\left.f \in\left(H^{2}(\Omega) \cap H_{0}^{1}(\Omega)\right)^{*}\right)$. The regularity of the solution $u$ depends on the given data $f$ and the domain geometry (Blum et al., 1980).

### 2.2 The naive mixed formulation

Intuitively, equation (1.1) can be decoupled into a system of two Poisson problems by introducing an auxiliary function $w$ such that

$$
\left\{\begin{array} { r l } 
{ - \Delta w = f } & { \text { in } \Omega , }  \tag{2.2}\\
{ w = 0 } & { \text { on } \partial \Omega , }
\end{array} \quad \text { and } \quad \left\{\begin{array}{rl}
-\Delta \bar{u}=w & \text { in } \Omega, \\
\bar{u}=0 & \text { on } \partial \Omega,
\end{array}\right.\right.
$$

where $(\bar{u}, w) \in H_{0}^{1}(\Omega) \times H_{0}^{1}(\Omega)$. We refer to (2.2) as the naive mixed formulation. Note that numerical solvers for the Poisson problems (2.2) are readily available, while numerical approximation of the fourth-order problem (1.1) is generally a much harder task. The mixed weak formulation of (2.2) is to find $\bar{u}, w \in H_{0}^{1}(\Omega)$ such that

$$
\begin{align*}
& A(w, \phi)=(f, \phi) \quad \forall \phi \in H_{0}^{1}(\Omega),  \tag{2.3a}\\
& A(\bar{u}, \psi)=(w, \psi) \quad \forall \psi \in H_{0}^{1}(\Omega), \tag{2.3b}
\end{align*}
$$

where

$$
A(\phi, \psi)=\int_{\Omega} \nabla \phi \cdot \nabla \psi \mathrm{d} x
$$

Given $f \in H^{-1}(\Omega) \subset\left(H^{2}(\Omega) \cap H_{0}^{1}(\Omega)\right)^{*}$, it is clear that the weak solutions $\bar{u}, w$ are well defined by (2.3) because they are solutions of decoupled Poisson problems (Evans, 1998). Since our goal is to solve the biharmonic problem (1.1), an important question is whether the solution $u$ in (2.1) and the solution $\bar{u}$ in (2.3) are the same.
Remark 2.1 For $f \in H^{-1}(\Omega)$, existing results suggest that under appropriate conditions, the solution $\bar{u}$ of system (2.3) is equivalent to the solution $u$ of equation (2.1) in the sense that

$$
u=\bar{u} \text { in } H^{2}(\Omega) \cap H_{0}^{1}(\Omega)
$$



Fig. 1. Domain $\Omega$ containing a reentrant corner.
These conditions include (i) the domain $\Omega$ and the given data $f$ being smooth, which can be verified by the regularity of these equations up to the domain boundary (Gilbarg \& Trudinger, 1983; Li \& Nistor, 2009), (ii) the polygonal domain $\Omega$ being convex (Zhang \& Zhang, 2008). It is however pointed out that $u$ is not always equivalent to $\bar{u}$ when the polygonal domain $\Omega$ has reentrant corners, which is known as the Sapongyan paradox (Nazarov \& Plamenevsky, 1994; Zhang \& Zhang, 2008). In this case, the numerical solution for (2.3) does not converge to the solution of the biharmonic problem (1.1). In the next subsection, we shall study the structure of the solution in the presence of a reentrant corner in order to design effective numerical algorithms for equation (1.1).

### 2.3 Image of the Laplace operator and its orthogonal space

From now on we assume the given function $f \in L^{2}(\Omega)$ in (1.1). In addition, assume that the polygonal domain $\Omega$ has a reentrant corner associated with the vertex $Q$ and the corresponding interior angle $\omega \in(\pi, 2 \pi)$. Without loss of generality, we set $Q$ to be the origin. Let $(r, \theta)$ be polar coordinates centered at the vertex $Q$, such that $\omega$ is spanned by two half-lines $\theta=0$ and $\theta=\omega$. Given $R>0$ we identify a sector $K_{\omega}^{R} \subset \Omega$ with radius $R$ as

$$
K_{\omega}^{R}=\{(r \cos \theta, r \sin \theta) \in \Omega \mid 0 \leq r \leq R, 0 \leq \theta \leq \omega\} .
$$

A sketch drawing of the domain $\Omega$ is shown in Fig. 1.
The mapping $-\Delta: H^{2}(\Omega) \cap H_{0}^{1}(\Omega) \rightarrow L^{2}(\Omega)$ is injective and has a closed range (Grisvard, 1992). Denote by $\mathcal{M}$ the image of this mapping and by $\mathcal{M}^{\perp}$ its orthogonal complement. Then it follows that $\mathcal{M} \oplus \mathcal{M}^{\perp}=L^{2}(\Omega)$. Therefore, if $w \in \mathcal{M}$ in (2.3), we have equivalent solutions $u=\bar{u}$ in $H^{2}(\Omega) \cap H_{0}^{1}(\Omega)$. When the domain $\Omega$ is convex, one has $\mathcal{M}=L^{2}(\Omega)$; namely, the solution of the Poisson equation with Dirichlet boundary condition is always in $H^{2}(\Omega)$ when $f \in L^{2}(\Omega)$. Thus, in a convex domain the condition $w \in H_{0}^{1}(\Omega) \subset \mathcal{M}$ holds, and therefore the solutions $u$ and $\bar{u}$ are equivalent. However, if $\Omega$ contains reentrant corners, $\mathcal{M}$ is a strict subset of $L^{2}(\Omega)$ and in general $w \notin \mathcal{M}$. Consequently, the solution in (2.2) $\bar{u} \notin H^{2}(\Omega)$ and it is different from the solution of (1.1) $u \in H^{2}(\Omega) \cap H_{0}^{1}(\Omega)$. Fortunately, the space $\mathcal{M}^{\perp}$ is finite-dimensional and it is possible to identify its basis.

We first introduce an $L^{2}$ function in domain $\Omega$ in the following way.
Definition 2.2 Given the parameters $\tau \in(0,1)$ and $R$ such that $K_{\omega}^{R} \subset \Omega$, we define an $L^{2}$ function in $\Omega$,

$$
\begin{equation*}
\xi(r, \theta ; \tau, R):=s^{-}(r, \theta ; \tau, R)+\zeta(r, \theta ; \tau, R), \tag{2.4}
\end{equation*}
$$

where

$$
\begin{equation*}
s^{-}(r, \theta ; \tau, R)=\eta(r ; \tau, R) r^{-\frac{\pi}{\omega}} \sin \left(\frac{\pi}{\omega} \theta\right) \in L^{2}(\Omega), \tag{2.5}
\end{equation*}
$$

with cut-off function $\eta(r ; \tau, R) \in C^{\infty}(\Omega)$ satisfying $\eta(r ; \tau, R)=1$ for $0 \leq r \leq \tau R$ and $\eta(r ; \tau, R)=0$ for $r>R$, and $\zeta \in H_{0}^{1}(\Omega)$ satisfies

$$
\begin{equation*}
-\Delta \zeta=\Delta s^{-} \text {in } \Omega, \quad \zeta=0 \text { on } \partial \Omega \tag{2.6}
\end{equation*}
$$

From (2.5) we see that $s^{-} \in C^{\infty}\left(\Omega \backslash K_{\omega}^{\delta}\right)$ for any $\delta>0$ and $s^{-}=0$ for $(r \cos \theta, r \sin \theta) \in \Omega \backslash K_{\omega}^{R}$. Moreover, $\Delta s^{-}=0$ if $r<\tau R$ or $r>R$.

Remark 2.3 The function $\xi$ and its variants have been used for solving the Poisson equation with singularities. For example, using a fixed coefficient $\tau=\frac{1}{2}$, $\xi$ was utilized in Blum \& Dobrowolski (1981), Cai \& Kim (2001), and another version of $\xi$ without the cut-off function was analyzed in Lions \& Magenes (1972), Ciarlet \& Raviart (1974). When solving the biharmonic problem (1.1), we shall see in the next section that all these versions of $\xi$ lead to similar finite element algorithms. However, the value of $\tau$ can affect the value of $s^{-}$in Algorithm 3.1, which will be elaborated on later.

Denote the $i$ th side of $\partial \Omega$ by $\bar{\Gamma}_{i}$, where $\Gamma_{i}$ is open. For the function $\xi$ defined in (2.4), we introduce the maximal extension of the Laplace operator in $L^{2}(\Omega)$ :

$$
D\left(\Delta, L^{2}(\Omega)\right):=\left\{v \in L^{2}(\Omega): \Delta v \in L^{2}(\Omega)\right\} .
$$

For a function $v$, let $\gamma_{i} v$ be the restriction of $v$ to $\Gamma_{i}$. Then the mapping

$$
v \mapsto \gamma_{i} v
$$

that is defined for $v \in H^{2}(\Omega)$ has a unique continuous extension from $D\left(\Delta, L^{2}(\Omega)\right)$ into $\widetilde{H}^{-\frac{1}{2}}\left(\Gamma_{i}\right)$ (Grisvard, 1992, Theorem 1.5.2), where $\widetilde{H}^{-\frac{1}{2}}\left(\Gamma_{i}\right)$ is the dual space of $\widetilde{H}^{\frac{1}{2}}\left(\Gamma_{i}\right)$ that is defined at the beginning of Section 2. Then we have the following properties of $\xi$.
Lemma 2.4 For a given $\eta \in C^{\infty}(\Omega)$ as defined in Definition 2.2, the function $\xi \in D\left(\Delta, L^{2}(\Omega)\right)$ is uniquely defined and satisfies

$$
\begin{equation*}
-\Delta \xi=0 \text { in } \Omega, \quad \xi=0 \text { on } \partial \Omega \tag{2.7}
\end{equation*}
$$

In addition, $\xi$ depends on the domain $\Omega$, but not on $\tau$ or $R$. Namely, for any $\tau_{1}, \tau_{2}$ and $R_{1}, R_{2}$ satisfying $0<\delta<\min \left\{\tau_{1} R_{1}, \tau_{2} R_{2}\right\}$, it follows that

$$
\begin{equation*}
\xi(r, \theta):=\xi\left(r, \theta ; \tau_{1}, R_{1}\right)=\xi\left(r, \theta ; \tau_{2}, R_{2}\right) . \tag{2.8}
\end{equation*}
$$

Proof. The condition in (2.7) follows from Grisvard (1992, Lemma 2.3.6). We proceed to prove (2.8). For $0<\delta<\min \left\{\tau_{1} R_{1}, \tau_{2} R_{2}\right\}$, we have $K_{\omega}^{\delta} \subset K_{\omega}^{\tau_{1} R_{1}} \cap K_{\omega}^{\tau_{2} R_{2}} \subset \Omega$. By (2.5) we have

$$
s^{-}\left(r, \theta ; \tau_{1}, R_{1}\right)-s^{-}\left(r, \theta ; \tau_{2}, R_{2}\right)=0, \quad(r, \theta) \in K_{\omega}^{\delta}
$$

Recall that $s^{-}\left(r, \theta ; \tau_{i}, R_{i}\right) \in C^{\infty}\left(\Omega \backslash K_{\omega}^{\delta}\right)$; then it follows that

$$
s^{-}\left(r, \theta ; \tau_{1}, R_{1}\right)-s^{-}\left(r, \theta ; \tau_{2}, R_{2}\right) \in C^{\infty}(\Omega) .
$$

Since $\zeta\left(r, \theta ; \tau_{i}, R_{i}\right) \in H_{0}^{1}(\Omega)$, we have

$$
\begin{aligned}
\tilde{\xi} & :=\xi\left(r, \theta ; \tau_{1}, R_{1}\right)-\xi\left(r, \theta ; \tau_{2}, R_{2}\right) \\
& =\zeta\left(r, \theta ; \tau_{1}, R_{1}\right)-\zeta\left(r, \theta ; \tau_{2}, R_{2}\right)+\left(s^{-}\left(r, \theta ; \tau_{1}, R_{1}\right)-s^{-}\left(r, \theta ; \tau_{2}, R_{2}\right)\right) \in H_{0}^{1}(\Omega)
\end{aligned}
$$

From (2.7) we have

$$
\Delta \tilde{\xi}=\Delta \xi\left(r, \theta ; \tau_{1}, R_{1}\right)-\Delta \xi\left(r, \theta ; \tau_{2}, R_{2}\right)=0 \text { in } \Omega, \quad \tilde{\xi}=0 \text { on } \partial \Omega .
$$

By the Lax-Milgram theorem, we have $\tilde{\xi}=0$, and thus (2.8) holds.
From now on we shall write $\xi(r, \theta)$ instead of $\xi(r, \theta ; \tau, R)$, since it is independent of $\tau$ and $R$. We also notice that $\xi(r, \theta) \not \equiv 0$, because otherwise we have $s^{-}=-\zeta \in H_{0}^{1}(\Omega)$, which contradicts the fact that $s^{-} \notin H_{0}^{1}(\Omega)$.
Remark 2.5 It is clear that $\xi \in D\left(\Delta, L^{2}(\Omega)\right)$ is an $L^{2}(\Omega)$ solution of the boundary value problem (2.7). It is interesting to note that different from the $H^{1}(\Omega)$ solution, the $L^{2}(\Omega)$ solution to problem (2.7) is not unique. For example, there are at least two $L^{2}(\Omega)$ solutions for (2.7), one is $\xi=0$, which is also the unique $H^{1}(\Omega)$ solution, and the other is $\xi \not \equiv 0$ as defined in (2.4). Note that the boundary condition in (2.7) is defined in the trace sense, namely $\xi \in \widetilde{H}^{-\frac{1}{2}}\left(\Gamma_{i}\right)$ on each side.

Now we are ready to describe the subspace $\mathcal{M}^{\perp}$. For the dimension of $\mathcal{M}^{\perp}$, we have the following results from Grisvard (1992, Theorem 2.3.7).

Lemma 2.6 The dimension of $\mathcal{M}^{\perp}$ is equal to the cardinality of the set $\left\{\lambda_{k}: 0<\lambda_{k}<1\right\}$ for $k \geq 1$, namely

$$
\operatorname{dim}\left(\mathcal{M}^{\perp}\right)=\operatorname{card}\left\{\lambda_{k}: 0<\lambda_{k}<1\right\}
$$

with $\lambda_{k}^{2}$ being the eigenvalues of the one-dimensional problem

$$
-\partial_{\theta \theta} \phi_{k}=\lambda_{k}^{2} \phi_{k} \quad \text { in }(0, \omega), \quad \phi(0)=\phi(\omega)=0 .
$$

For $k \geq 1$, it is clear that when $\lambda_{k}>0$,

$$
\begin{equation*}
\lambda_{k}=\frac{k \pi}{\omega}, \quad \phi_{k}=\sqrt{\frac{2}{\omega}} \sin \left(\frac{k \pi}{\omega} \theta\right) . \tag{2.9}
\end{equation*}
$$

Hence, for the domain $\Omega$ with one reentrant corner, $\mathcal{M}^{\perp}$ satisfies the following lemma.
Lemma 2.7 The dimension of $\mathcal{M}^{\perp}$ is $\operatorname{dim}\left(\mathcal{M}^{\perp}\right)=1$ and $\mathcal{M}^{\perp}=\operatorname{span}\{\xi(r, \theta)\}$, where $\xi(r, \theta)$ is the $L^{2}$ function defined in (2.4).

By Lemma 2.7, any $w \in L^{2}(\Omega)$ can be uniquely expressed as

$$
w=w_{\mathcal{M}}+c \xi
$$

where $w_{\mathcal{M}}=w-c \xi \in \mathcal{M}$ and the coefficient

$$
\begin{equation*}
c=\frac{(w, \xi)-\left(w_{\mathcal{M}}, \xi\right)}{\|\xi\|^{2}}=\frac{(w, \xi)}{\|\xi\|^{2}} \tag{2.10}
\end{equation*}
$$

### 2.4 The modified mixed formulation

Based on the discussion above, we propose a modified mixed formulation for (1.1),

$$
\left\{\begin{array} { r l } 
{ - \Delta w } & { = f \quad \text { in } \Omega , }  \tag{2.11}\\
{ w } & { = 0 \quad \text { on } \partial \Omega , }
\end{array} \quad \text { and } \quad \left\{\begin{array}{rl}
-\Delta \tilde{u} & =w-c \xi \quad \text { in } \Omega \\
\tilde{u} & =0 \quad \text { on } \partial \Omega
\end{array}\right.\right.
$$

where $\xi$ is given in (2.4) and the coefficient $c$ is shown in (2.10). The corresponding modified mixed weak formulation for (2.11) is to find $\tilde{u}, w \in H_{0}^{1}(\Omega)$ such that

$$
\begin{align*}
& A(w, \phi)=(f, \phi),  \tag{2.12a}\\
& A(\tilde{u}, \psi)=(w-c \xi, \psi), \tag{2.12b}
\end{align*}
$$

for any $\phi, \psi \in H_{0}^{1}(\Omega)$.
Then we have the following result for the modified mixed formulation.
Theorem 2.8 Given $f \in L^{2}(\Omega)$, let $\tilde{u}$ be the solution of the modified mixed weak formulation (2.12) and let $u$ be the solution of the weak formulation (2.1). Then it follows that $\tilde{u} \in H^{2}(\Omega) \cap H_{0}^{1}(\Omega)$ and $u=\tilde{u}$.

Proof. Since $f \in L^{2}(\Omega)$ we have $w \in H_{0}^{1}(\Omega) \subset L^{2}(\Omega)$. Thus it follows that $w-c \xi \in \mathcal{M}$, which implies $\tilde{u} \in H^{2}(\Omega) \cap H_{0}^{1}(\Omega)$. Then (2.12b) becomes

$$
\begin{equation*}
-(\Delta \tilde{u}, \psi)=(w-c \xi, \psi) \quad \forall \psi \in H_{0}^{1}(\Omega) . \tag{2.13}
\end{equation*}
$$

Note $\Delta \tilde{u} \in L^{2}(\Omega)$. Then following the density argument, (2.13) leads to

$$
-(\Delta \tilde{u}, \psi)=(w-c \xi, \psi), \quad \forall \psi \in L^{2}(\Omega)
$$

Thus, for any $\phi \in H^{2}(\Omega) \cap H_{0}^{1}(\Omega)$, we have $\Delta \phi \in L^{2}(\Omega)$ and therefore

$$
\begin{equation*}
(\Delta \tilde{u}, \Delta \phi)=(w-c \xi,-\Delta \phi) . \tag{2.14}
\end{equation*}
$$

Since $\xi \in L^{2}(\Omega), \Delta \xi=0 \in L^{2}(\Omega)$, using Green's theorem (Grisvard, 1992, Theorem 1.5.3) and (2.7), we have that $(\xi, \Delta \phi)=0$. Then the right-hand side of (2.14) becomes

$$
(w-c \xi,-\Delta \phi)=A(w, \phi)+(c \xi, \Delta \phi)=A(w, \phi)=(f, \phi),
$$

where the last equation is based on (2.12a). Hence we have obtained that $\tilde{u} \in H^{2}(\Omega) \cap H_{0}^{1}(\Omega)$ satisfies

$$
(\Delta \tilde{u}, \Delta \phi)=(f, \phi) \quad \forall \phi \in H^{2}(\Omega) \cap H_{0}^{1}(\Omega) .
$$

This is the same equation as (2.1) that defines $u$. Consequently, $\tilde{u}=u \in H^{2}(\Omega) \cap H_{0}^{1}(\Omega)$ and we have completed the proof.

Therefore, by Theorem 2.8 , the solution $u$ of the biharmonic problem (1.1) satisfies

$$
\left\{\begin{array} { r l } 
{ - \Delta w } & { = f \quad \text { in } \Omega , }  \tag{2.15}\\
{ w } & { = 0 \quad \text { on } \partial \Omega , }
\end{array} \quad \text { and } \quad \left\{\begin{array}{rl}
-\Delta u & =w-c \xi \quad \text { in } \Omega, \\
u & =0 \quad \text { on } \partial \Omega
\end{array}\right.\right.
$$

The corresponding weak formulation is to find $u, w \in H_{0}^{1}(\Omega)$ such that for any $\phi, \psi \in H_{0}^{1}(\Omega)$,

$$
\begin{align*}
& A(w, \phi)=(f, \phi),  \tag{2.16a}\\
& A(u, \psi)=(w-c \xi, \psi), \tag{2.16b}
\end{align*}
$$

where $c$ is given in (2.10).
In addition, we have the following regularity result.
Lemma 2.9 Given $f \in L^{2}(\Omega)$, for $w, u$ in the modified mixed formulation (2.15), it follows that

$$
\begin{align*}
\|w\|_{H^{1}(\Omega)} & \leq C\|f\|  \tag{2.17a}\\
\|u\|_{H^{2}(\Omega)} & \leq C\|f\| . \tag{2.17b}
\end{align*}
$$

Proof. Estimate (2.17a) is a direct consequence of the fact that the Laplace operator is an isomorphism between $H_{0}^{1}(\Omega)$ and $H^{-1}(\Omega)$ and $\|f\|_{H^{-1}(\Omega)} \leq C\|f\|$. By the Lax-Milgram theorem for (2.1), we have that $\Delta$ is also an isomorphism from $H^{2}(\Omega) \cap H_{0}^{1}(\Omega)$ to its dual. Therefore, estimate (2.17b) follows from this observation and Theorem 2.8.

## 3. The finite element method

In this section we propose a linear $C^{0}$ finite element method solving the biharmonic problem (1.1). Then we derive the finite element error analysis for the solution $u$ to show that our method will achieve the optimal convergence rate especially when the domain is nonconvex.

### 3.1 The finite element algorithm

Let $\mathcal{T}_{n}$ be a triangulation of $\Omega$ with shape-regular triangles and let $S_{n} \subset H_{0}^{1}(\Omega)$ be the $C^{0}$ Lagrange finite element space associated with $\mathcal{T}_{n}$,

$$
\begin{equation*}
S_{n}(\mathcal{T}):=\left\{v \in C^{0}(\Omega) \cap H_{0}^{1}(\Omega):\left.v\right|_{T} \in P_{1} \forall T \in \mathcal{T}_{n}\right\} \tag{3.1}
\end{equation*}
$$

where $P_{1}$ is the space of polynomials of degree no more than 1 . Then we proceed to propose the finite element algorithm.

Algorithm 3.1 We define the finite element solution of the biharmonic problem (1.1) by utilizing the decoupling in (2.16) as follows.

- Step 1. Find the finite element solution $w_{n} \in S_{n}$ of the Poisson equation

$$
\begin{equation*}
A\left(w_{n}, \phi\right)=(f, \phi) \quad \forall \phi \in S_{n} \tag{3.2}
\end{equation*}
$$

- Step 2. With $s^{-}$defined in (2.5), we compute the finite element solution $\zeta_{n} \in S_{n}$ of the Poisson equation

$$
\begin{equation*}
A\left(\zeta_{n}, \phi\right)=\left(\Delta s^{-}, \phi\right) \quad \forall \phi \in S_{n} \tag{3.3}
\end{equation*}
$$

and set $\xi_{n}=\zeta_{n}+s^{-}$.

- Step 3. Find the coefficient $c_{n} \in \mathbb{R}$ such that

$$
\int_{\Omega}\left(w_{n}-c_{n} \xi_{n}\right) \xi_{n} \mathrm{~d} x=0
$$

or equivalently, we compute the coefficient

$$
\begin{equation*}
c_{n}=\frac{\left(w_{n}, \xi_{n}\right)}{\left\|\xi_{n}\right\|^{2}}=\frac{\left(w_{n}, \zeta_{n}\right)+\left(w_{n}, s^{-}\right)}{\left\|s^{-}\right\|^{2}+\left\|\zeta_{n}\right\|^{2}+2\left(\zeta_{n}, s^{-}\right)} \tag{3.4}
\end{equation*}
$$

- Step 4. Find the finite element solution $u_{n} \in S_{n}$ of the Poisson equation

$$
\begin{equation*}
A\left(u_{n}, \psi\right)=\left(w_{n}-c_{n} \xi_{n}, \psi\right) \quad \forall \psi \in S_{n} \tag{3.5}
\end{equation*}
$$

Remark 3.2 The function $s^{-}$in (2.5) exists only in the presence of a reentrant corner. When the domain is convex, we set $s^{-}=0$ and Algorithm 3.1 reduces to the naive mixed finite element algorithm for equation (1.1). According to (3.3), $\zeta_{n} \in S_{n}$, while $\xi_{n} \in L^{2}(\Omega)$ but $\xi_{n} \notin S_{n}$. In addition, the finite element approximations in Algorithm 3.1 are well defined based on the Lax-Milgram theorem.

Remark 3.3 All the integrals in Algorithm 3.1 can be effectively approximated by quadrature rules except for the term $\left\|s^{-}\right\|^{2}=\int_{K_{\omega}^{R}}\left(s^{-}\right)^{2} \mathrm{~d} x$ in (3.4), due to the lack of regularity in the integrand. Namely, $s^{-} \in H^{\beta}\left(K_{\omega}^{R}\right)$ with $\beta<1-\frac{\pi}{\omega}$. Only evaluating $\left\|s^{-}\right\|^{2}$ by quadrature rules may destroy the convergence rate of the proposed algorithm. In practice, we evaluate $\left\|s^{-}\right\|^{2}$ in the following way:

$$
\begin{align*}
\left\|s^{-}\right\|^{2} & =\int_{K_{\omega}^{R}}\left(s^{-}\right)^{2} \mathrm{~d} x=\int_{K_{\omega}^{\tau R}}\left(s^{-}\right)^{2} \mathrm{~d} x+\int_{K_{\omega}^{R} \backslash K_{\omega}^{\tau R}}\left(s^{-}\right)^{2} \mathrm{~d} x \\
& =\int_{0}^{\omega}\left(\int_{0}^{\tau R}+\int_{\tau R}^{R}\right)\left(\eta(r ; \tau, R) r^{-\frac{\pi}{\omega}} \sin \left(\frac{\pi}{\omega} \theta\right)\right)^{2} r \mathrm{~d} r \mathrm{~d} \theta \\
& =\int_{0}^{\omega} \int_{0}^{\tau R} r^{1-\frac{2 \pi}{\omega}} \sin ^{2}\left(\frac{\pi}{\omega} \theta\right) \mathrm{d} r \mathrm{~d} \theta+\int_{0}^{\omega} \int_{\tau R}^{R} \eta^{2}(r ; \tau, R) r^{1-\frac{2 \pi}{\omega}} \sin ^{2}\left(\frac{\pi}{\omega} \theta\right) \mathrm{d} r \mathrm{~d} \theta  \tag{3.6}\\
& =\frac{\omega(\tau R)^{2-\frac{2 \pi}{\omega}}}{4-\frac{4 \pi}{\omega}}+\frac{\omega}{2} \int_{\tau R}^{R} \eta^{2}(r ; \tau, R) r^{1-\frac{2 \pi}{\omega}} \mathrm{~d} r
\end{align*}
$$

where the one-dimensional definite integral in the last equation can be calculated directly or evaluated by a quadrature rule in one dimension. The computational technique (3.6) is accurate, and at the same time can significantly reduce the computational cost in evaluating the integral.

### 3.2 Optimal error estimates on quasi-uniform meshes

Suppose that the mesh $\mathcal{T}_{n}$ consists of quasi-uniform triangles with size $h$. Recall the interpolation error estimate on $\mathcal{T}_{n}$ (Ciarlet Jr \& He, 2003) for any $v \in H^{l}(\Omega), l>1$ :

$$
\begin{equation*}
\left\|v-v_{I}\right\|_{H^{m}(\Omega)} \leq C h^{l-m}\|v\|_{H^{l}(\Omega)} \tag{3.7}
\end{equation*}
$$

where $m=0,1$ and $v_{I} \in S_{n}$ represents the nodal interpolation of $v$. For the Poisson equations (2.6) and (2.15) in the polygonal domain with a reentrant corner, given $f \in L^{2}(\Omega)$, it is well known that $w, \zeta \in H^{\alpha}(\Omega)$ with $\alpha<1+\frac{\pi}{\omega}$ (see e.g., Grisvard, 1985, 1992). Recall the finite element approximations $w_{n}$ and $\zeta_{n}$ in (3.2) and (3.3), respectively. Due to the lack of regularity, the standard error estimate (Ciarlet Jr \& He, 2003) yields

$$
\begin{equation*}
\left\|w-w_{n}\right\|_{H^{1}(\Omega)} \leq C h^{\alpha-1}\|w\|_{H^{\alpha}(\Omega)}, \quad\left\|\zeta-\zeta_{n}\right\|_{H^{1}(\Omega)} \leq C h^{\alpha-1}\|\zeta\|_{H^{\alpha}(\Omega)} \tag{3.8}
\end{equation*}
$$

Note that $\xi-\xi_{n}=\zeta-\zeta_{n} \in H^{\alpha}(\Omega)$, and thus

$$
\begin{equation*}
\left\|\xi-\xi_{n}\right\|_{H^{1}(\Omega)} \leq C h^{\alpha-1}\|\zeta\|_{H^{\alpha}(\Omega)} \tag{3.9}
\end{equation*}
$$

In addition, we have the following $L^{2}$ error analysis.
Lemma 3.4 Given $w_{n}$ and $\xi_{n}$ in Algorithm 3.1, we have

$$
\begin{equation*}
\left\|w-w_{n}\right\| \leq C h^{2 \alpha-2}\|w\|_{H^{\alpha}(\Omega)}, \quad\left\|\xi-\xi_{n}\right\| \leq C h^{2 \alpha-2}\|\zeta\|_{H^{\alpha}(\Omega)} \tag{3.10}
\end{equation*}
$$

Proof. We only prove the error estimate for $w-w_{n}$, and the estimate for $\xi-\xi_{n}\left(=\zeta-\zeta_{n}\right)$ can be obtained similarly. Consider the Poisson problem

$$
\begin{equation*}
-\Delta v=g \text { in } \Omega, \quad v=0 \text { on } \partial \Omega, \tag{3.11}
\end{equation*}
$$

where $g \in L^{2}(\Omega)$. By the Aubin-Nitsche lemma in Ciarlet Jr \& He (2003, Theorem 3.2.4), we have

$$
\begin{equation*}
\left\|w-w_{n}\right\| \leq C\left\|w-w_{n}\right\|_{H^{1}(\Omega)} \sup _{g \neq 0 \in L^{2}(\Omega)}\left(\frac{\inf _{\phi \in S_{n}}\|v-\phi\|_{H^{1}(\Omega)}}{\|g\|}\right) . \tag{3.12}
\end{equation*}
$$

For $\alpha \in\left(1,1+\frac{\pi}{\omega}\right) \subset(1,2)$, the regularity result (Grisvard, 1985, 1992) gives $\|v\|_{H^{\alpha}(\Omega)} \leq$ $C_{\alpha}\|g\|_{H^{\alpha-2}} \leq C\|g\|$. Then we have

$$
\begin{equation*}
\inf _{\phi \in S_{n}}\|v-\phi\|_{H^{1}(\Omega)} \leq\left\|v-v_{I}\right\|_{H^{1}(\Omega)} \leq C h^{\alpha-1}\|v\|_{H^{\alpha}(\Omega)} \leq C h^{\alpha-1}\|g\| . \tag{3.13}
\end{equation*}
$$

Combining the estimates in (3.13), (3.12) and (3.8), we have completed the proof.

Next we carry out the error estimate for the finite element approximation $u_{n}$ in (3.5).
Theorem 3.5 Let $u_{n} \in S_{n}$ be the finite element approximation to (3.5), and $u$ be the solution to the biharmonic problem (2.1). Then it follows that

$$
\left\|u-u_{n}\right\|_{H^{1}(\Omega)} \leq C h
$$

Proof. For any $\phi \in S_{n}$, based on (2.16b) and (3.5), we have

$$
\begin{aligned}
A(u, \phi) & =(w, \phi)-c(\xi, \phi) \\
A\left(u_{n}, \phi\right) & =\left(w_{n}, \phi\right)-c_{n}\left(\xi_{n}, \phi\right) .
\end{aligned}
$$

Taking the difference of the two equations above, we have

$$
\begin{align*}
A\left(u-u_{n}, \phi\right) & =\left(w-w_{n}, \phi\right)+c_{n}\left(\xi_{n}, \phi\right)-c(\xi, \phi)  \tag{3.14}\\
& =\left(w-w_{n}, \phi\right)+c_{n}\left(\xi_{n}-\xi, \phi\right)+\left(c_{n}-c\right)(\xi, \phi)
\end{align*}
$$

Let $u_{I} \in S_{n}$ be the nodal interpolation of $u$. Set $\epsilon=u_{I}-u, e=u_{I}-u_{n}$ and take $\phi=e$ in (3.14). We have

$$
A(e, e)=A(\epsilon, e)+\left(w-w_{n}, e\right)+c_{n}\left(\xi_{n}-\xi, e\right)+\left(c_{n}-c\right)(\xi, e) .
$$

Thus, we have

$$
\|e\|_{H^{1}(\Omega)}^{2} \leq C\left(\|\epsilon\|_{H^{1}(\Omega)}+\left\|w-w_{n}\right\|_{H^{-1}(\Omega)}+\left|c_{n}\right|\left\|\xi_{n}-\xi\right\|_{H^{-1}(\Omega)}+\left|c-c_{n}\right|\|\xi\|_{H^{-1}(\Omega)}\right)\|e\|_{H^{1}(\Omega)}
$$

Using the triangle inequality and the inequality above we have

$$
\begin{align*}
\left\|u-u_{n}\right\|_{H^{1}(\Omega)} & \leq\|e\|_{H^{1}(\Omega)}+\|\epsilon\|_{H^{1}(\Omega)} \\
& \leq C\left(\|\epsilon\|_{H^{1}(\Omega)}+\left\|w-w_{n}\right\|_{H^{-1}(\Omega)}+\left|c_{n}\right|\left\|\xi_{n}-\xi\right\|_{H^{-1}(\Omega)}+\left|c-c_{n}\right|\|\xi\|_{H^{-1}(\Omega)}\right) \\
& \leq C\left(\|\epsilon\|_{H^{1}(\Omega)}+\left\|w-w_{n}\right\|+\left|c_{n}\right|\left\|\xi_{n}-\xi\right\|+\left|c-c_{n}\right|\|\xi\|\right) \tag{3.15}
\end{align*}
$$

The last inequality is based on the fact that the $H^{-1}$ norm of an $L^{2}$ function is bounded by its $L^{2}$ norm. We shall estimate every term in (3.15). Recall the solution $u \in H^{2}(\Omega)$. By the interpolation error estimate (3.7),

$$
\begin{equation*}
\|\epsilon\|_{H^{1}(\Omega)}=\left\|u-u_{I}\right\|_{H^{1}(\Omega)} \leq C h\|u\|_{H^{2}(\Omega)} . \tag{3.16}
\end{equation*}
$$

Recall the angle of the reentrant corner $\omega \in(\pi, 2 \pi)$. Thus, choosing $\alpha=3 / 2<1+\frac{\pi}{\omega}$ in (3.10), we have

$$
\begin{equation*}
\left\|w-w_{n}\right\| \leq C h, \quad\left\|\xi-\xi_{n}\right\| \leq C h \tag{3.17}
\end{equation*}
$$

To obtain the error estimate for the third term in (3.15), we still need to show that $\left|c_{n}\right|$ is uniformly bounded. Recall that $\xi \not \equiv 0$ depends only on the domain $\Omega$, and so is

$$
\begin{equation*}
\|\xi\|>0 \tag{3.18}
\end{equation*}
$$

Moreover, when $h \leq h_{0}:=\min \left\{1, \sqrt[2 \alpha-2]{\frac{\|\xi\|}{2 C\|\xi\|_{H^{\alpha}(\Omega)}}}\right\}$, it follows from (3.10) that

$$
\begin{equation*}
\frac{1}{2}\|\xi\| \leq\left\|\xi_{n}\right\| \leq \frac{3}{2}\|\xi\| . \tag{3.19}
\end{equation*}
$$

By setting $\phi=w_{n}$ in (3.2) and applying the Poincaré inequality, we obtain

$$
\begin{equation*}
\left\|w_{n}\right\| \leq C\left\|w_{n}\right\|_{H^{1}(\Omega)} \leq C\|f\| \tag{3.20}
\end{equation*}
$$

By (3.4), (3.19) and (3.20), we have the uniform boundedness,

$$
\begin{equation*}
\left|c_{n}\right| \leq \frac{\left\|w_{n}\right\|}{\left\|\xi_{n}\right\|} \leq C\|f\| \tag{3.21}
\end{equation*}
$$

where $C$ is a constant depending on $\Omega$. Subtracting (3.4) from (2.10) we obtain

$$
c-c_{n}=\frac{\left(w-w_{n}, \xi\right)}{\|\xi\|^{2}}+\frac{\left(\xi-\xi_{n}, w_{n}\right)}{\left\|\xi_{n}\right\|^{2}}+\frac{\left\|\xi_{n}\right\|^{2}-\|\xi\|^{2}}{\|\xi\|^{2}\left\|\xi_{n}\right\|^{2}}\left(w_{n}, \xi\right) .
$$

By the uniform boundedness for terms in (3.18), (3.19), (3.20) and the estimates in (3.17), we have for the last term in (3.15),

$$
\begin{equation*}
\left|c-c_{n}\right|\|\xi\| \leq\left\|w-w_{n}\right\|+\frac{\left(\left\|\xi_{n}\right\|+2\|\xi\|\right)\left\|w_{n}\right\|}{\left\|\xi_{n}\right\|^{2}}\left\|\xi-\xi_{n}\right\| \leq C h . \tag{3.22}
\end{equation*}
$$

Then the proof is completed by plugging (3.16), (3.21), (3.22) and (3.17) into (3.15).
Remark 3.6 The error estimate in Theorem 3.5 shows that the proposed finite element algorithm (Algorithm 3.1) produces numerical solutions that converge to the solution of the biharmonic problem (1.1) when the domain $\Omega$ is nonconvex. In the case that $\Omega$ is convex, Algorithm 3.1 reduces to the naive mixed finite element algorithm for equation (1.1) that has proven to be effective (Zhang \& Zhang, 2008). Therefore, Algorithm 3.1 approximates the target equation in both convex and nonconvex domains. On a quasi-uniform mesh, the convergence is first order (optimal) for $u$ in the $H^{1}$ norm (Theorem 3.5) and suboptimal (3.8) for the auxiliary function $w$ in the $H^{1}$ norm. In Algorithm 3.1, one shall solve three Poisson problems. Given the availability of fast Poisson solvers, Algorithm 3.1 is a relatively easy and cost effective alternative to existing algorithms solving (1.1).

## 4. Optimal error estimates on graded meshes

The numerical approximations from Algorithm 3.1 are optimal for $u$ but only sub-optimal for $w$. It is largely due to the lack of regularity for $w$. Recall that for $f \in L^{2}(\Omega), w$ is merely in $H^{\alpha}(\Omega)$ for
$\alpha<1+\frac{\pi}{\omega}$. In this section we study system (2.16) in a class of weighted Sobolev spaces and in turn propose graded triangulations that lead to numerical solutions converging at the optimal rate to both $u$ and $w$.

### 4.1 Regularity in weighted Sobolev spaces

We now introduce the Kondratiev-type weighted spaces for the analysis of system (2.16).
Definition 4.1 [Weighted Sobolev spaces] Recall that $Q$ is the vertex at the reentrant corner. Let $r(x)$ be the distance from $x$ to $Q$. For $a \in \mathbb{R}, m \geq 0$ and $G \subset \Omega$, we define the weighted Sobolev space

$$
\mathcal{K}_{a}^{m}(G):=\left\{v, r^{|\alpha|-a} \partial^{\alpha} v \in L^{2}(G) \forall|\alpha| \leq m\right\},
$$

where the multi-index $\alpha=\left(\alpha_{1}, \alpha_{2}\right) \in \mathbb{Z}_{\geq 0}^{2},|\alpha|=\alpha_{1}+\alpha_{2}$ and $\partial^{\alpha}=\partial_{x}^{\alpha_{1}} \partial_{y}^{\alpha_{2}}$. The $\mathcal{K}_{a}^{m}(G)$ norm for $v$ is defined by

$$
\|v\|_{\mathcal{K}_{a}^{m}(G)}=\left(\sum_{|\alpha| \leq m} \int_{G}\left|r^{|\alpha|-a} \partial^{\alpha} v\right|^{2} \mathrm{~d} x\right)^{\frac{1}{2}}
$$

Remark 4.2 According to Definition 4.1, in the region that is away from the reentrant corner, the weighted space $\mathcal{K}_{a}^{m}$ is equivalent to the Sobolev space $H^{m}$. In the neighborhood of $Q$, the space $\mathcal{K}_{a}^{m}$ is the same Kondratiev space (Kondrat'ev, 1967; Grisvard, 1985; Dauge, 1988). Recall the first equation in (2.15) that defines $w$. In the Dirichlet Poisson problem, the reentrant corner can give rise to singularities in $w$, such that $w \notin H^{2}(\Omega)$. It is the reason that the finite element approximation to $w$ on a quasi-uniform mesh is not optimal. The singularity in $w$ is however local and concentrates in the neighborhood of $Q$. Involving a proper weight function, the space $\mathcal{K}_{a}^{m}$ may allow more singular functions and is an important tool for analyzing corner singularities.

In the weighted Sobolev space we have the following regularity result for system (2.16).
Lemma 4.3 Assume $a<\frac{\pi}{\omega}$ and $f \in L^{2}(\Omega)$. Recall $\zeta$ in (2.6). Then it follows that

$$
\|\zeta\|_{\mathcal{K}_{a+1}^{2}(\Omega)} \leq C\left\|\Delta s^{-}\right\|
$$

In addition, recall $w$ in (2.15). Then we have

$$
\|w\|_{\mathcal{K}_{a+1}^{2}(\Omega)} \leq C\|f\|
$$

Proof. Since $\Delta s^{-}, f \in L^{2}(\Omega) \subset \mathcal{K}_{a-1}^{0}(\Omega)$, the desired estimates follow by applying Li \& Nicaise (2018, Theorem 3.3) to equations (2.6) and (2.15).

### 4.2 Graded meshes

We now present the construction of graded meshes to improve the convergence rate of the numerical approximation from Algorithm 3.1.

Algorithm 4.4 [Graded refinements] Let $\mathcal{T}$ be a triangulation of $\Omega$ with shape-regular triangles. Recall that $Q$ is the vertex of $\Omega$ at the reentrant corner. It is clear that $Q$ is also a vertex in the triangulation


Fig. 2. The new node on an edge $p q$ (left - right): $p \neq Q$ and $q \neq Q$ (midpoint); $p=Q(|p r|=\kappa|p q|, \kappa<0.5)$.


Fig. 3. Refinement of a triangle $\Delta x_{0} x_{1} x_{2}$. First row: (left to right): the initial triangle and the midpoint refinement; second row: two consecutive graded refinements toward $x_{0}=Q,(\kappa<0.5)$.
$\mathcal{T}$. Let $p q$ be an edge in the triangulation $\mathcal{T}$ with $p$ and $q$ as the endpoints. Then, in a graded refinement, a new node $r$ on $p q$ is produced according to the following conditions:

1. (Neither $p$ or $q$ coincides with $Q$.) We choose $r$ as the midpoint $(|p r|=|q r|)$.
2. ( $p$ coincides with $Q$.) We choose $r$ such that $|p r|=\kappa|p q|$, where $\kappa \in(0,0.5)$ is a parameter that will be specified later. See Fig. 2 for example.
Then the graded refinement, denoted by $\kappa(\mathcal{T})$, proceeds as follows. For each triangle $T \in \mathcal{T}$, a new node is generated on each edge of as described above. Then $T$ is decomposed into four small triangles by connecting these new nodes (Fig. 3). Given an initial mesh $\mathcal{T}_{0}$ satisfying the condition above, the associated family of graded meshes $\left\{\mathcal{T}_{n}, n \geq 0\right\}$ is defined recursively $\mathcal{T}_{n+1}=\kappa\left(\mathcal{T}_{n}\right)$.

Given a grading parameter $\kappa$, Algorithm 4.4 produces smaller elements near $Q$ for better approximation of a singular solution. It is an explicit construction of graded meshes based on recursive refinements. See also Apel et al. (1996), Băcuţă et al. (2005), Krendl et al. (2016), Li \& Nicaise (2018) and the references therein for more discussions on the graded mesh. Note that after $n$ refinements, the number of triangles in the mesh $\mathcal{T}_{n}$ is $\mathcal{O}\left(4^{n}\right)$.

### 4.3 Optimal error estimates on graded meshes

In the rest of this section we shall show that with a proper selection of the grading parameter $\kappa$, the proposed numerical solutions $u_{n}$ and $w_{n}$ converge to the solutions $u$ and $w$ of (2.15) at the optimal rate on graded meshes. Recall the finite element space $S_{n}$ in (3.1) associated with the graded mesh $\mathcal{T}_{n}$.

We first recall the following interpolation error estimates (Apel et al., 1996; Li \& Nicaise, 2018) for functions in the weighted space.
Lemma 4.5 Let $0<a<\frac{\pi}{\omega}$ and choose the grading parameter $\kappa=2^{-1 / a}$. Define $h:=2^{-n}$. Then for any $v \in \mathcal{K}_{a+1}^{2}(\Omega)$, it follows that

$$
\left\|v-v_{I}\right\|_{H^{1}(\Omega)} \leq C h\|v\|_{\mathcal{K}_{a+1}^{2}(\Omega)}
$$

where $v_{I}$ is the nodal interpolation of $v$ associated with $\mathcal{T}_{n}$.
Now we proceed to derive the optimal error estimates of $w-w_{n}$ and $\xi-\xi_{n}$ on graded meshes.
Lemma 4.6 Let $0<a<\frac{\pi}{\omega}$ and choose $\kappa=2^{-1 / a}$. Then for the approximations $w_{n}$ and $\xi_{n}$ defined in (3.2) and (3.3), it follows that

$$
\begin{aligned}
&\left\|w-w_{n}\right\|_{H^{1}(\Omega)} \leq C h\|w\|_{\mathcal{K}_{a+1}^{2}(\Omega)}, \quad\left\|w-w_{n}\right\| \leq C h^{2}\|w\|_{\mathcal{K}_{a+1}^{2}(\Omega)} \\
&\left\|\xi-\xi_{n}\right\|_{H^{1}(\Omega)} \leq C h\|\zeta\|_{\mathcal{K}_{a+1}^{2}(\Omega)}, \quad\left\|\xi-\xi_{n}\right\| \leq C h^{2}\|\zeta\|_{\mathcal{K}_{a+1}^{2}(\Omega)}
\end{aligned}
$$

where $h:=2^{-n}$.
Proof. The $H^{1}$ error estimates for $w-w_{n}$ and $\zeta-\zeta_{n}$ follow from the standard $H^{1}$ error estimates on graded meshes (Apel et al., 1996). We only prove the $L^{2}$ error estimates for $w-w_{n}$ and the estimates for $\zeta-\zeta_{n}$ will follow similarly. Applying the Aubin-Nitsche lemma to (3.11) again, we have

$$
\begin{equation*}
\left\|w-w_{n}\right\| \leq C\left\|w-w_{n}\right\|_{H^{1}(\Omega)} \sup _{g \neq 0 \in L^{2}(\Omega)}\left(\frac{\inf _{\phi \in S_{n}}\|v-\phi\|_{H^{1}(\Omega)}}{\|g\|}\right) . \tag{4.1}
\end{equation*}
$$

Based on the regularity estimate in Lemma 4.3, the function $v$ in (3.11) satisfies

$$
\|v\|_{\mathcal{K}_{a+1}^{2}(\Omega)} \leq C\|g\|
$$

Together with Lemma 4.5, we have

$$
\begin{equation*}
\inf _{\phi \in S_{n}}\|v-\phi\|_{H^{1}(\Omega)} \leq\left\|v-v_{I}\right\|_{H^{1}(\Omega)} \leq C h\|v\|_{\mathcal{K}_{a+1}^{2}(\Omega)} \leq C h\|g\|, \tag{4.2}
\end{equation*}
$$

where $v_{I}$ is the nodal interpolation of $v$ associated with $\mathcal{T}_{n}$. Plugging (4.2) into (4.1) leads to the desired error estimate for $\left\|w-w_{n}\right\|$.

The $H^{1}$ and $L^{2}$ error estimates of $\xi-\xi_{n}$ follow by the relationship $\xi-\xi_{n}=\zeta-\zeta_{n}$.

Remark 4.7 Note that the error estimates for $w-w_{n}, \zeta-\zeta_{n}$ follow from the standard $H^{1}$ error estimate on both uniform meshes (Ciarlet Jr \& He, 2003) and graded meshes (Apel et al., 1996). However, the $H^{1}$ and $L^{2}$ error estimates of $\xi-\xi_{n}$ are obtained based on the relationship $\xi-\xi_{n}=\zeta-\zeta_{n}$.

We conclude this section with the $H^{1}$ error estimate for the solution $u$ of the biharmonic problem (1.1) on graded meshes.

Theorem 4.8 Let $\mathcal{T}_{n}$ be the graded mesh with the grading parameter $0<\kappa<\frac{1}{2}$. Let $u_{n}$ be the finite element approximation to $u$ that is defined in Algorithm 3.1. Then it follows that

$$
\left\|u-u_{n}\right\|_{H^{1}(\Omega)} \leq C h,
$$

where $h:=2^{-n}$.

Proof. Let $u_{I}$ be the nodal interpolation of $u$ associated with $\mathcal{T}_{n}$. Similar to the analysis in Theorem 3.5 on quasi-uniform meshes, we have

$$
\begin{equation*}
\left\|u-u_{n}\right\|_{H^{1}(\Omega)} \leq C\left(\left\|u-u_{I}\right\|_{H^{1}(\Omega)}+\left\|w-w_{n}\right\|+\frac{\left\|w_{n}\right\|}{\left\|\xi_{n}\right\|}\left\|\xi_{n}-\xi\right\|+\left|c-c_{n}\right|\|\xi\|_{H^{-1}(\Omega)}\right) . \tag{4.3}
\end{equation*}
$$

For $0<\kappa<\frac{1}{2}$, the following interpolation error still holds:

$$
\begin{equation*}
\left\|u-u_{I}\right\|_{H^{1}(\Omega)} \leq C h\|u\|_{H^{2}(\Omega)} . \tag{4.4}
\end{equation*}
$$

Thus, the proof is completed by combining the estimates in (4.3), (4.4), (3.22) and the $L^{2}$ error estimates for $w-w_{n}$ and $\xi-\xi_{n}$ in (3.10).

Remark 4.9 According to Theorems 3.5 and 4.8, the numerical solution $u_{n}$ in Algorithm 3.1 approximates the solution $u$ of the biharmonic problem in the optimal $H^{1}$ convergence rate on quasiuniform meshes and also on graded meshes defined in Algorithm 4.4. In addition, a proper graded mesh can improve the effectiveness in approximating the auxiliary function $w$ in (2.15). In particular, selecting the grading parameter $\kappa$ as in Lemma 4.6, the proposed finite element solution $w_{n}$ converges to $w$ in both $H^{1}$ and $L^{2}$ norms with the optimal rate on graded meshes. Nonetheless, the numerical approximations $u_{n}$ and $w_{n}$ from Algorithm 3.1 converge to $u$ and $w$ in both convex and nonconvex domains. The graded mesh can improve the convergence rate but does not make divergent numerical solutions convergent. We also point out that when high-order finite element methods are used in Algorithm 3.1, new graded meshes are needed to recover the optimal $H^{1}$ convergence rate for both $u$ and $w$. We shall study these cases in future works.

## 5. Numerical illustrations

In this section we present numerical test results to validate our theoretical predictions for the proposed finite element method solving equation (1.1). Since the solutions $u, w$ in (2.15) are unknown, we use the numerical convergence rate

$$
\begin{equation*}
\mathcal{R}=\log _{2} \frac{\left|v_{j}-v_{j-1}\right|_{H^{1}(\Omega)}}{\left|v_{j+1}-v_{j}\right|_{H^{1}(\Omega)}} \tag{5.1}
\end{equation*}
$$



FIg. 4. The square domain (Example 5.1): (a) the initial mesh; (b) the mesh after one refinement; (c): the solution $u_{7}$ from Algorithm 3.1; (d) $\left|u_{R}-u_{7}\right|$.

Table $1 \quad H^{1}$ Convergence history of the $P_{1}$ elements in Example 5.1 on uniform meshes

|  | $j=3$ | $j=4$ | $j=5$ | $j=6$ |
| :--- | :---: | :---: | :---: | :---: |
| $\mathcal{R}$ for $u_{j}$ | 0.96 | 0.99 | 1.00 | 1.00 |
| $\mathcal{R}$ for $w_{j}$ | 0.96 | 0.99 | 1.00 | 1.00 |

as an indicator of the actual convergence rate. Here $v_{j}$ denotes the finite element solution on the mesh $\mathcal{T}_{j}$ obtained after $j$ refinements of the initial triangulation $\mathcal{T}_{0}$. It can be either $u_{j}$ or $w_{j}$ depending on the underlying Poisson problem. In particular, suppose the actual convergence rate is $\left\|v-v_{j}\right\|_{H^{1}(\Omega)}=\mathcal{O}\left(h^{\beta}\right)$ for $\beta>0$. Then for the $P_{1}$ finite element method, the rate in (5.1) is a good approximation of the exponent $\beta$ as the level of refinements $j$ increases (Li et al., 2010).

We shall use the solution of the $C^{0}$ interior penalty discontinuous Galerkin ( $C^{0}$-IPDG) method (Brenner, 2012) programmed in the software FEniCS (Alnæs et al., 2015) as a reference solution. More specifically, we use the $C^{0}$-IPDG method based on $P_{2}$ polynomials with the penalty parameter $\eta=24$. The reference solution is computed on the mesh after seven mesh refinements of the given initial mesh and is denoted by $u_{R}$. Since the $C^{0}$-IPDG method leads to numerical solutions converging to the solution $u$ regardless of the convexity of the domain, we can use $u_{R}$ as a good approximation of $u$.

Example 5.1 (A convex domain). We consider problem (1.1) with $f=10$ in the square domain $\Omega=(0,2)^{2}$. Since all the vertices have angles less than $\pi$, Algorithm 3.1 coincides with the naive mixed finite element method based on the formulations in (2.3).

We solve this problem using Algorithm 3.1 on uniform meshes obtained by midpoint refinements with the initial mesh given in Fig. 4(a). The finite element solution $u_{7}$ and the difference $\left|u_{R}-u_{7}\right|$ are shown in Figs 4(c) and 4(d), respectively. The convergence rates (5.1) for $u_{j}$ and $w_{j}$ on a sequence of uniform meshes are shown in Table 1. We see that the solution of the mixed finite element method converges to the solution of the biharmonic equation (1.1) and the optimal convergence rate $(\mathcal{R}=1)$ is achieved for both the numerical solution $u_{j}$ and the auxiliary finite element solution $w_{j}$. This is consistent with our expectation (Remark 2.1) for the problem in a convex domain.

Example 5.2 (A nonconvex domain). In this example, we investigate the convergence of Algorithm 3.1 by considering equation (1.1) with $f=1$ in an L-shaped domain $\Omega=\Omega_{0} \backslash \Omega_{1}$ with $\Omega_{0}=(-2,2)^{2}$


FIG. 5. The L-shaped domain (Example 5.2): (a) the initial mesh; (b) the graded mesh after two refinements; (c) the solution $u_{7}^{U}$ of the naive mixed method; (d) the difference $\left|u_{R}-u_{7}^{U}\right|$; (e) the solution $u_{7}^{A}$ from Algorithm 3.1; (f) the difference $\left|u_{R}-u_{7}^{A}\right|$.
and $\Omega_{1}=(0,2) \times(-2,0)$. We use the following cut-off function in the algorithm:

$$
\eta(r ; \tau, R)= \begin{cases}0 & \text { if } r \geq R, \\ 1 & \text { if } r \leq \tau R, \\ \frac{1}{2}-\frac{15}{16}\left(\frac{2 r}{R(1-\tau)}-\frac{1+\tau}{1-\tau}\right)+\frac{5}{8}\left(\frac{2 r}{R(1-\tau)}-\frac{1+\tau}{1-\tau}\right)^{3}-\frac{3}{16}\left(\frac{2 r}{R(1-\tau)}-\frac{1+\tau}{1-\tau}\right)^{5} & \text { otherwise },\end{cases}
$$

where $R=\frac{9}{5}, \tau=\frac{1}{8}$.
In the first test, we solve equation (1.1) in the L-shaped domain using quasi-uniform meshes and compare the performances of Algorithm 3.1 and the naive mixed finite element algorithm based on formulation (2.3). On $\mathcal{T}_{j}$ we denote the numerical solutions from Algorithm 3.1 by $u_{j}^{A}$ and $w_{j}^{A}$, and denote the numerical solutions from the naive mixed finite element algorithm by $u_{j}^{U}$ and $w_{j}^{U}$. The initial mesh is shown in Fig. 5(a). In Table 2 we display the errors ( $u_{R}-u_{j}^{U}$ and $u_{R}-u_{j}^{A}$ ) in the $L^{\infty}$ norm between the finite element solutions and the reference solution $u_{R}$. In addition, the differences $\left|u_{R}-u_{7}^{U}\right|$ and $\left|u_{R}-u_{7}^{A}\right|$ are presented in Figs 5(d) and 5(f). From these results we see that the solution $u_{j}^{A}$ from Algorithm 3.1 converges to the actual solution, while the solution $u_{j}^{U}$ of the naive mixed finite element algorithm does not converge to the solution of the biharmonic equation (1.1) as the meshes are refined. These

Table 2 The $L^{\infty}$ error in the L-shaped domain on quasi-uniform meshes

|  | $j=3$ | $j=4$ | $j=5$ | $j=6$ |
| :--- | :---: | :---: | :---: | :---: |
| $\left\\|u_{R}-u_{j}^{U}\right\\|_{L^{\infty}(\Omega)}$ | $1.28014 \mathrm{e}-01$ | $1.37318 \mathrm{e}-01$ | $1.41309 \mathrm{e}-01$ | $1.42525 \mathrm{e}-01$ |
| $\left\\|u_{R}-u_{j}^{A}\right\\|_{L^{\infty}(\Omega)}$ | $1.58074 \mathrm{e}-02$ | $7.84320 \mathrm{e}-03$ | $3.20391 \mathrm{e}-03$ | $1.20794 \mathrm{e}-03$ |

Table 3 Numerical convergence rates $\mathcal{R}$ for $u_{j}^{U}$ and $w_{j}^{U}$ in the L-shaped domain

| $\kappa \backslash j$ | $u_{j}^{U}$ |  |  |  |  | $w_{j}^{U}$ |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $j=5$ | $j=6$ | $j=7$ | $j=8$ | $j=9$ | $j=5$ | $j=6$ | $j=7$ | $j=8$ | $j=9$ |
| $\kappa=0.1$ | 0.95 | 0.98 | 0.99 | 1.00 | 1.00 | 0.95 | 0.98 | 0.99 | 1.00 | 1.00 |
| $\kappa=0.2$ | 0.96 | 0.99 | 1.00 | 1.00 | 1.00 | 0.96 | 0.99 | 0.99 | 1.00 | 1.00 |
| $\kappa=0.3$ | 0.97 | 0.98 | 0.99 | 0.99 | 0.99 | 0.96 | 0.98 | 0.99 | 0.99 | 0.99 |
| $\kappa=0.4$ | 0.94 | 0.94 | 0.94 | 0.93 | 0.93 | 0.94 | 0.95 | 0.95 | 0.94 | 0.94 |
| $\kappa=0.5$ | 0.86 | 0.82 | 0.78 | 0.75 | 0.72 | 0.87 | 0.84 | 0.80 | 0.77 | 0.74 |

observations are closely aligned with our theoretical predictions in Remark 3.6. Namely, Algorithm 3.1 gives rise to convergent numerical solutions in both convex and nonconvex domains, while the naive mixed method is applicable only for convex domains.

For the naive mixed finite element method, although the solution $u_{j}^{U}$ does not converge to the solution of (1.1), we notice that both $u_{j}^{U}$ and $w_{j}^{U}, j \geq 0$ are converging sequences. The numerical convergence rate $\mathcal{R}$ (5.1) for $u_{j}^{U}$ and $w_{j}^{U}$ on a sequence of graded meshes (including quasi-uniform meshes) is reported in Table 3. In the table we observe that $u_{j}^{U}$ and $w_{j}^{U}$ have similar convergence rates: $\mathcal{R}<1$ on quasi-uniform meshes and on the graded meshes with $\kappa=0.4$; and $\mathcal{R}=1$ on graded meshes with $\kappa \leq 0.3$. These results indicate that the naive mixed finite element solution $u_{j}^{U}$ converges to the solution $\bar{u}$ of (2.2) in $H^{1}(\Omega)$. Recall however that $\bar{u} \neq u$ when the domain has reentrant corners.

In the last test, we examine the convergence rates of the finite element solution $u_{j}^{A}$ and the auxiliary finite element solution $w_{j}^{A}$ from Algorithm 3.1 on a sequence of graded meshes (including quasi-uniform meshes). The $H^{1}$ convergence rates (5.1) for the finite element solutions $u_{j}^{A}$ and $w_{j}^{A}$ are reported in Table 4. For $u_{j}^{A}$, the optimal convergence rate $(\mathcal{R}=1)$ is achieved on all meshes with $\kappa \in(0,0.5]$. For the auxiliary solution $w_{j}^{A}$, we observe that the convergence rate is not optimal on quasi-uniform meshes and on the graded meshes with $\kappa=0.4$; and the optimal convergence rate $\mathcal{R}=1$ is obtained on graded meshes when $\kappa \leq 0.3$. These numerical results justify the theory (Theorems 3.5, 4.8 and Lemma 4.6) developed early in this paper. Namely, the numerical solution $u_{j}^{A}$ converges to $u$ at the optimal rate on quasi-uniform meshes and on graded meshes, while $w_{j}^{A}$ will converge to $w$ in the optimal rate when $\kappa<2^{-\frac{\omega}{\pi}}=2^{-\frac{3}{2}} \approx 0.354$. For $\kappa>2^{-\frac{3}{2}}\left(\kappa=0.4,0.5\right.$ in Table 4), $w_{j}^{A}$ shall converge to $w$ at a reduced rate due to the fact that $w$ is singular near the reentrant corner $\left(w \notin H^{2}(\Omega)\right.$ and $w \in H^{\alpha}(\Omega)$ for $\alpha<1+\frac{2}{3} \approx 1.667$ see; (3.8)).
Example 5.3 In this example we compare the CPU time and the memory usage of the proposed finite element algorithm (Algorithm 3.1) with those of the $C^{0}$-IPDG method programmed in FEniCS. More specifically, the tests are based on these two algorithms/methods for biharmonic problems in

Table $4 \quad H^{1}$ convergence history in the L-shaped domain on graded meshes

| $\kappa \backslash j$ | $u_{j}^{A}$ |  |  |  |  | $w_{j}^{A}$ |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $j=5$ | $j=6$ | $j=7$ | $j=8$ | $j=9$ | $j=5$ | $j=6$ | $j=7$ | $j=8$ | $j=9$ |
| $\kappa=0.1$ | 0.95 | 0.98 | 0.99 | 1.00 | 1.00 | 0.95 | 0.98 | 0.99 | 1.00 | 1.00 |
| $\kappa=0.2$ | 0.95 | 0.99 | 1.00 | 1.00 | 1.00 | 0.96 | 0.99 | 0.99 | 1.00 | 1.00 |
| $\kappa=0.3$ | 0.97 | 0.99 | 1.00 | 1.00 | 1.00 | 0.96 | 0.98 | 0.99 | 0.99 | 0.99 |
| $\kappa=0.4$ | 0.98 | 0.99 | 1.00 | 1.00 | 1.00 | 0.94 | 0.95 | 0.95 | 0.94 | 0.94 |
| $\kappa=0.5$ | 0.97 | 0.99 | 0.99 | 1.00 | 1.00 | 0.87 | 0.84 | 0.80 | 0.77 | 0.74 |

Table 5 The CPU time (in seconds) of Algorithm 3.1 and the $C^{0}$-IPDG method in FEniCS. ('-' represents running out of memory)

| method $\backslash j$ | Example 5.1 |  |  |  |  | Example 5.2 |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $j=5$ | $j=6$ | $j=7$ | $j=8$ | $j=9$ | $j=5$ | $j=6$ | $j=7$ | $j=8$ | $j=9$ |
| $C^{0}$-IPDG | 0.65 | 4.38 | 31.85 | 279.19 | - | 0.56 | 3.30 | 24.12 | 189.15 | - |
| Algorithm 3.1 | 0.07 | 0.28 | 1.21 | 5.08 | 24.2 | 0.12 | 0.46 | 1.82 | 8.15 | 35.60 |

Table 6 The memory usage (in GB) of Algorithm 3.1 and the $C^{0}$-IPDG method in FEniCS. ('-' represents running out of memory)

| method $\backslash j$ | Example 5.1 |  |  |  |  | Example 5.2 |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $j=5$ | $j=6$ | $j=7$ | $j=8$ | $j=9$ | $j=5$ | $j=6$ | $j=7$ | $j=8$ | $j=9$ |
| $C^{0}$-IPDG | 0.14 | 0.34 | 1.28 | 5.74 | - | 0.12 | 0.25 | 0.91 | 4.04 | - |
| Algorithm 3.1 | $<0.01$ | 0.01 | 0.02 | 0.28 | 1.26 | $<0.01$ | $<0.01$ | 0.01 | 0.16 | 0.94 |

both Examples 5.1 and 5.2 on quasi-uniform meshes. The results of the CPU time (in seconds) of both algorithms/methods are shown in Table 5. Here, the CPU time does not include the mesh generation time since we assume they are the same. The results of the memory usage (in GB) are shown in Table 6. In this example, all results are tested on Ubuntu 20.04 with 16 GB memory and Intel Core ${ }^{\mathrm{TM}}$ i7-6600U processors.

From Table 5 we find that given the same triangulation, the implementation of the proposed finite element algorithm (Algorithm 3.1) can be much faster than that of the $C^{0}$-IPDG method, due to the availability of fast Poisson solvers. At the same time, Table 6 indicates that Algorithm 3.1 uses less memory compared with the $C^{0}$-IPDG method. These results demonstrate the efficacy of our algorithm for solving the biharmonic problem (1.1).

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## Conflict of interest

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